

New Approach for Numerical Solution of Poisson's Equation by Cubic Spline

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Received December 23, 2014; Revised December 29, 2014; Accepted December 31, 2014

Abstract We consider the solution of various boundary value problems for Poisson's equation in the unit square using a nodal cubic spline collocation method and modifications of it which produce optimal fourth order approximations. Uniform partition of the square with cost $O(N^2 \log N)$ using a direct fast Fourier transform method. The numerical results exhibit super convergence phenomena.

Keywords: Nodal collocation, Poisson's equation, cubic spline, convergence analysis

Cite This Article: M. Yousefi, J. Rashidinia, M. Yousefi, N.S BahrololoumiMofrad, and Mehdi Moudi, "New Approach for Numerical Solution of Poisson's Equation by Cubic Spline." *Journal of Mathematical Sciences and Applications*, vol. 2, no. 3 (2014): 39-42. doi: 10.12691/jmsa-2-3-3.

1. Introduction

The first spline collocation method proposed for the solution of second-order two-point boundary value problems was a nodal cubic spline collocation (NCSC) method [1]. In the basic nodal cubic spline collocation (NCSC) method, an approximate solution is sought in the space of bicubic splines and is determined by satisfying the differential equation and the boundary conditions at the partition nodes of Ω and $\partial\Omega$, respectively. This method is well known to be suboptimal; in fact, it is no better than second order, whereas fourth order is expected since the approximate solution is piecewise bicubic [2]. In [3], two optimal order NCSC methods, a two step method (TSM) and a one step method (OSM), were presented for the solution of the Dirichlet BVP. These methods are based on optimal NCSC methods for solving second-order two-point BVPs. De Boor [4] proved that classical nodal cubic spline collocation for solving two-point boundary value problems. Archer [5] and, independently, Daniel and Swartz [6] developed a modified nodal cubic spline collocation (MNCSC) scheme which is fourth order accurate. Houstis, Vavalis, and Rice [7]. Derived a fourth order MNCSC scheme for solving elliptic boundary value problems on rectangles. In [5], for the Helmholtz equation, matrix decomposition algorithms (MDA) with fast Fourier transforms were formulated and implemented to solve the two step method (TSM) collocation equations for each the boundary value problem (BVP), but, for the one step method, it was possible to formulate an MDA only for the Dirichlet BVP.

In this paper, we consider Poisson's equation with a constant coefficient in the unit square subject to the Dirichlet boundary value condition. We develop a new algorithm for the solution of the linear systems arising in the NCSC method for the Dirichlet problem. Our scheme involves perturbations of both the left- and right-hand sides. Numerical results show that our scheme exhibits super convergence phenomena while that [7] does not.

A brief outline of this paper is as follows. We give preliminaries in section 2. In section 3, the matrix vector of our scheme, and a direct fast Fourier transforms algorithm for solving the scheme are presented Section 4.) Includes numerical results obtained using our scheme.

2. Cubic Spline Function Properties and Cubic Spline Interpolation

We consider the Dirichlet boundary value problem for Poisson's equation

$$\Delta u = f \text{ in } \Omega, \quad u = 0 \text{ on } \partial\Omega \quad (2.1)$$

Where $\Delta \equiv D_x^2 + D_y^2$ is the Laplacian, and $\Omega = (0,1) \times (0,1)$, and $\partial\Omega$ is the boundary of Ω .

Let $\rho_x = \{x_i\}_{i=0}^{N+1}$ with $x_i = ih, h = 1/(N+1)$ be a uniform partition of $[0,1]$ in the x-direction (For simplicity, in the remainder of the paper, a uniform partition

$\rho_y = \{y_i\}_{i=0}^{N+1}$ of $[0,1]$ in the y-direction is such that $y_j = x_j$.)

Let S_3 denote the space of cubic splines

$$S_3 = \{v : v \in C^2([0,1]), v|_I \in P_3, 1 \leq i \leq N+1\}$$

Where $I_i = [x_{i-1}, x_i]$ and P_3 is the set of polynomials of degree ≤ 3 and let

$$S^D = \{v \in S_3 : v(0) = v(1) = 0\}$$

To introduce basis functions for S_3 , we extend the partition ρ_x using $N+2, N+3, N+4$. As a basis for S_3 , we choose the functions $\{\beta_m\}_{m=-1}^{N+2}$, where

$$\beta_m = \begin{cases} h^{-3}g_1(x-x_{m-2}), & x \in I_{m-1} \\ g_2((x-x_{m-1})/h), & x \in I_m \\ g_2((x_{m+1}-x)/h), & x \in I_{m+1} \\ h^{-3}g_1(x_{m+2}-x), & x \in I_{m+2} \\ 0 & \text{otherwise} \end{cases} \quad (2.2)$$

And

$$g_1(x) = x^3, \quad g_2(x) = 1 + 3x + 3x^2 - 3x^3 \quad (2.3)$$

These basis functions are such that, for $m = 0, \dots, N+1$

$$\begin{cases} \beta_{m-1}(x_m) = 1, \beta_m(x_m) = 4, \beta_{m+1}(x_m) = 1, \\ \beta_{m-1}''(x_m) = 6/h^2, \beta_m''(x_m) = -12/h^2, \beta_{m+1}''(x_m) = 6/h^2. \end{cases} \quad (2.4)$$

As a basis for S^D , we choose the cubic splines $\{\beta_m^D\}_{m=0}^{N+1}$ defined in terms of splines by

$$\begin{cases} \beta_0^D = \beta_0 - 4\beta_{-1}, \beta_1^D = \beta_1 - \beta_{-1}, \beta_m^D = \beta_m, m = 2, 3, \dots, N-1 \\ \beta_N^D = \beta_N - \beta_{N+2}, \beta_{N+1}^D = \beta_{N+1} - \beta_{N+2} \end{cases} \quad (2.5)$$

Cf.[8,Section 2].

It follows from (2.3) that

$$\begin{cases} B_0^D(x_1) = 1, & B_1^D(x_1) = 4, & B_1^D(x_2) = 1, \\ B_N^D(x_{N-1}) = 1, & B_N^D(x_N) = 4, & B_{N+1}^D(x_N) = 1, \end{cases} \quad (2.6)$$

$$\begin{cases} [B_0^D]''(x_0) = -\frac{36}{h^2}, & [B_1^D]''(x_0) = 0, \\ [B_0^D]''(x_1) = \frac{6}{h^2}, & [B_1^D]''(x_1) = -\frac{12}{h^2}, & [B_1^D]''(x_2) = \frac{6}{h^2} \\ [B_N^D]''(x_{N-1}) = \frac{6}{h^2}, & [B_N^D]''(x_N) = -\frac{12}{h^2}, & [B_{N+1}^D]''(x_N) = \frac{6}{h^2} \\ [B_N^D]''(x_{N+1}) = \frac{6}{h^2}, & [B_{N+1}^D]''(x_{N+1}) = -\frac{36}{h^2}. \end{cases} \quad (2.7)$$

It also follows from (2.5),(2.6),(2.4), and (2.3) that, for $i=1, \dots, N$,

$$B_m^D(x_i) - \frac{h^2}{6}[B_m^D]''(x_i) = \begin{cases} 6, m=i \\ 0, m \neq i \end{cases} \quad m = 0, \dots, N+1 \quad (2.8)$$

And the matrix-vector form of

$$\phi_{i,j} = \sum_{m=1}^N c_{i,m} \sum_{n=1}^N c_{j,n} \psi_{m,n}, \quad i, j = 1, \dots, N \quad (2.9)$$

Is

$$\phi = (C_1 \otimes C_2)\psi \quad (2.10)$$

Where $C_1 = (c_{i,m}^{(1)})_{i \in I, m \in M}$, $C_2 = (c_{j,n}^{(2)})_{j \in J, n \in N}$. And

$$\begin{aligned} \phi &= [\phi_{1,1}, \dots, \phi_{1,N}, \dots, \phi_{N,1}, \dots, \phi_{N,N}]^T, \\ \psi &= [\psi_{1,1}, \dots, \psi_{1,N}, \dots, \psi_{N,1}, \dots, \psi_{N,N}]^T \end{aligned}$$

3. The Approximate Solution of Poisson's Equation

In this section, we consider equation (2.1). we introduce

$$S^D = S_3 \cap \{v : v(0) = v(1) = 0\}$$

Our NCSC scheme for solving (2.1) is formulated as follows: seek

$u_h^D \in S^D \otimes S^D$ such that

$$\begin{aligned} \Delta u_h(x_i, y_j) - \frac{h^2}{36} D_x^2 D_y^2 u_h(x_i, y_j) \\ = f(x_i, y_j) - \frac{h^2}{144} \Delta f(x_i, y_j) \quad (3.1) \\ i, j = 0, \dots, N+1 \\ u = 0 \quad \text{on } \partial\Omega \end{aligned}$$

(3.1) is equivalent to

$$2D_x^2 D_y^2 u_h(x_i, y_j) = \Delta f(x_i, y_j), \quad i, j = 0, N+1 \quad (3.2)$$

$$\begin{aligned} D_x^2 u_h(x_i, y_j) - \frac{h^2}{36} D_x^2 D_y^2 u_h(x_i, y_j) \\ = f(x_i, y_j) - \frac{h^2}{144} \Delta f(x_i, y_j) \quad i = 0, N+1, j = 1, \dots, N \end{aligned} \quad (3.3)$$

$$\begin{aligned} D_y^2 u_h(x_i, y_j) - \frac{h^2}{36} D_x^2 D_y^2 u_h(x_i, y_j) \\ = f(x_i, y_j) - \frac{h^2}{144} \Delta f(x_i, y_j) \quad (3.4) \\ i = 1, \dots, N, j = 0, N+1 \end{aligned}$$

$$\begin{aligned} \Delta u_h(x_i, y_j) - \frac{h^2}{36} D_x^2 D_y^2 u_h(x_i, y_j) \\ = f(x_i, y_j) - \frac{h^2}{144} \Delta f(x_i, y_j) \quad i, j = 1, \dots, N \end{aligned} \quad (3.5)$$

The scheme (3.1) is motivated by the fourth order finite difference method for (2.1).

For $i = 0, N+1, 1 \leq j \leq N$ and $1 \leq i \leq N, 0 \leq j \leq N+1$, and

$$D_y^l D_x^2 u_h^D(x_i, y_j) = D_y^l f(x_i, y_j), \quad i, j = 0, N+1 \quad (3.6)$$

Where l in (3.6) is either 1 or 2.
 Since

$$u_h^D(x, y) = \sum_{m=0}^{N+1} \sum_{n=0}^{N+1} u_{m,n} \beta_m^D(x) \beta_n^D(y) \quad (3.7)$$

And involve $(N+2)^2$ equations in the unknown coefficients $\{u_{m,n}\}_{m,n=0}^{N+1}$

It follows from (2.4) and (2.5) that, for $0 \leq m \leq N+1$

$$2 \sum_{m=0}^{N+1} \sum_{n=0}^{N+1} u_{m,n} [B_m^D]''(x_i) [B_n^D]''(y_j) = \Delta f(x_i, y_j) \quad (3.8)$$

Substituting (3.7) into (3.2), we obtain

$$\sum_{m=0}^{N+1} \sum_{n=0}^{N+1} u_{m,n} \left(\begin{aligned} & [B_m^D]''(x_i) B_n^D(y_j) \\ & + \left[B_m^D(x_i) - \frac{h^2}{36} [B_n^D]''(x_j) \right] [B_n^D]''(y_j) \end{aligned} \right) \quad (3.9)$$

$$= f(x_i, y_j) - \frac{h^2}{144} \Delta f(x_i, y_j) \quad i, j = 1, \dots, N$$

Substituting (3.7) into, we obtain (3.3)

$$\sum_{m=1}^N \sum_{n=1}^N u_{m,n} \left(\begin{aligned} & [B_m^D]''(x_i) B_n^D(y_j) \\ & + \left[B_m^D(x_i) - \frac{h^2}{36} [B_n^D]''(x_j) \right] [B_n^D]''(y_j) \end{aligned} \right) \quad (3.10)$$

$$= \omega_{i,j} \quad i, j = 1, \dots, N$$

Where

$$\omega_{i,j} = f(x_i, y_j) - \frac{h^2}{144} \Delta f(x_i, y_j)$$

$$- \sum_{m=0, N+1} \sum_{n=0}^{N+1} u_{m,n} \left(\begin{aligned} & [B_m^D]''(x_i) B_n^D(y_j) \\ & + \left[B_m^D(x_i) - \frac{h^2}{36} [B_n^D]''(x_j) \right] [B_n^D]''(y_j) \end{aligned} \right)$$

$$- \sum_{m=1}^N \sum_{n=0, N+1} u_{m,n} \left(\begin{aligned} & [B_m^D]''(x_i) B_n^D(y_j) \\ & + \left[B_m^D(x_i) - \frac{h^2}{36} [B_n^D]''(x_j) \right] [B_n^D]''(y_j) \end{aligned} \right)$$

Each of the methods gives rise to a linear system of the form

$$A_1 \otimes B_2 + (B_1 \otimes A_2)u = f$$

Where the matrices A_1, B_1 are $M \times M$ with $M = N, N+1, \text{ or } N+2$, where the matrices A_2, B_2 are $(N+2) \times (N+2)$, and where \otimes denotes the matrix tensor product (cf., [8]), thus we can (3.10) as

$$\left[A \otimes B + \left(B - \frac{h^2}{36} A \right) \otimes A \right] u = v \quad (3.11)$$

Where $u = [u_{1,1}, \dots, u_{N,N}]^T, \lambda = [\lambda_{1,1}, \dots, \lambda_{N,N}]^T$ and for $\{B_m^D\}_{m=0}^{N+1}$, we introduce $N \times N$ matrices A and B defined by

$$A = (a_{i,m})_{i,m=1}^N, a_{i,m} = [B_m^D]_{j,n=1}^N, \quad (3.12)$$

$$B = (b_{j,n})_{j,n=1}^N b_{j,n} = B_n^D(y_j)$$

It follows from (2.4)-(2.7) that

$$A = 6h^{-2}P, \quad B = P + 6I \quad (3.13)$$

Where I is the identity matrix and the $N \times N$ matrix P is given by

$$P = \begin{bmatrix} -2 & 1 & & & \\ 1 & -2 & 1 & & \\ & \ddots & \ddots & \ddots & \\ & & & 1 & -2 & 1 \\ & & & & & 1 & -2 \end{bmatrix} \quad (3.14)$$

And (3.11) simplifies to

$$\frac{36}{h^2} [36P \otimes I + (36I + P) \otimes P] u = v \quad (3.15)$$

We have

$$K^{-1}PK = \Lambda \quad (3.16)$$

Where the $N \times N$ matrices Λ and Q are given by

$$\Lambda = \text{diag}(\lambda_i)_{i=1}^N, \quad \lambda_i = -4 \sin^2 \frac{i\pi}{2(N+1)} \quad (3.17)$$

$$K = (k_{i,j})_{i,j=1}^N, \quad k_{i,j} = \left(\frac{2}{N+1} \right)^{1/2} \sin \frac{ij\pi}{N+1} \quad (3.18)$$

Cf., [1, section 2].

Using (3.18), we see (3.15) is equivalent to

$$36h^{-2} (K \otimes I) [36P \otimes I + (36I + P) \otimes P] \quad (3.19)$$

$$(K \otimes I) (K^{-1} \otimes I) u = (K \otimes I) v$$

If $u' = (K^{-1} \otimes I)u, v' = (K \otimes I)v$ and using (3.16) and (3.19) we obtain

$$36h^{-2} [36\Lambda \otimes I + (36I + \Lambda) \otimes T] u' = v' \quad (3.20)$$

The system (3.20) reduces to the N independent systems

$$36h^{-2} [36\lambda_i \otimes I + (36I + \lambda_i) \otimes T] u'_i = v'_i \quad (3.21)$$

We thus have the following algorithm for solving (3.15) :

- Step 1. Compute $v' = (K \otimes I)v$
- Step 2. Solve the N systems in (3.21)
- Step 3. Compute $u = (K \otimes I)u'$

It follows from (3.18) that the cost of steps 1 and 3 is $O(N^2 \log N)$ each (cf.,[8]).In step 2, the systems are tridiagonal, so this step is performed at a cost $O(N^2)$. Consequently, the total cost of the algorithm is $O(N^2 \log N)$.

$$\|u\|_h = \max_{0 \leq i, j \leq N+1} |u(x_i, y_j)|, \|u\|_{C(\bar{\Omega})} \approx \max_{0 \leq i, j \leq 1001} |u(x_i, y_j)|$$

Where

$$t_i = \frac{i}{1001}, \quad i = 1, \dots, 1001.$$

Convergence rates were determines using the formula

$$rate = \frac{\log(e_{N/2} / e_N)}{\log[(N + 1) / (N / 2 + 1)]}$$

Where e_N is the error corresponding to partition $\rho_x \times \rho_y$.

Problem D-1: Poisson's equation with exact solution

$$u(x, y) = 3e^{xy}(x^2 - x)(y^2 - y)$$

4. Numerical Experiment

In our numerical study, we used the following testproblem. Thecomputations were carried out in double precision. We determined the nodal and global errors using the formulas. test problem from [7] were considered.

We determined the nodal and global errors using the formulas

Table 1. Nodal errors and convergence rates for u, u_x, u_y

N	$\ u - u_h\ _h$		$\ (u - u_h)_x\ _h$		$\ (u - u_h)_y\ _h$		$\ (u - u_h)_{xy}\ _h$	
	Error	Rate	Error	Rate	Error	Rate	Error	Rate
8	6.014 - 06		8.234 - 05		7.008 - 05		1.897 - 03	
16	4.999 - 07	4.087	7.001 - 06	4.002	6.015 - 06	4.027	1.569 - 04	3.987
32	3.326 - 09	4.036	4.879 - 07	4.051	5.298 - 07	4.002	1.231 - 05	4.001
64	1.007 - 10	4.001	3.451 - 08	4.034	3.876 - 08	4.001	1.204 - 06	3.861
128	0.246 - 11	3.986	2.002 - 09	4.000	2.068 - 09	3.997	1.001 - 07	4.007
256								

Table 2. Global errors and convergence rates for u, u_x, u_y and u_{xy}

N	$\ u - u_h\ _{C(\bar{\Omega})}$		$\ (u - u_h)_x\ _{C(\bar{\Omega})}$		$\ (u - u_h)_y\ _{C(\bar{\Omega})}$		$\ (u - u_h)_{xy}\ _{C(\bar{\Omega})}$	
	Error	Rate	Error	Rate	Error	Rate	Error	Rate
8	7.799 - 06		0.987 - 04		0.999 - 04		1.985 - 03	
16	6.013 - 07	3.769	1.082 - 05	3.994	0.854 - 05	3.945	2.001 - 04	3.938
32	5.239 - 08	3.989	1.120 - 06	4.002	0.832 - 06	4.008	2.087 - 05	3.874
64	4.001 - 09	4.001	1.148 - 07	4.081	0.901 - 07	3.756	2.732 - 06	3.646
128	2.017 - 10	4.003	1.278 - 08	3.896	0.965 - 08	4.001	2.865 - 07	4.000
256								

5. Conclusion

We see from the results in Table 1 and Table 2 that Scheme produces fourth order accuracy for u in both the discrete and the continuous maximum norms. We also observe super convergence phenomena since the derivative approximation at the partition nodes are of order four.

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